

Lesson #8: Computational Methods for Plasma Physics Hines ①

I. Simulation of Single Particle Motion

A. Particle Motion in Specified Fields $\underline{E}(x,t)$ and $\underline{B}(x,t)$

1. Lorentz Force Law $\frac{d\underline{v}(t)}{dt} = \frac{q}{m} [\underline{E}(x,t) + \underline{v}(t) \times \underline{B}(x,t)]$

2. Velocity $\frac{d\underline{x}(t)}{dt} = \underline{v}(t)$

3. This is a set of Ordinary differential equations (ODEs) which can be solved, if $\underline{E}(x,t)$ & $\underline{B}(x,t)$ are known, to advance $\underline{x}(t)$ and $\underline{v}(t)$ in time.

4. This requires initial conditions at $t=0$, $\underline{x}(t=0)$ and $\underline{v}(t=0)$.

5. In general, any ODE problem can be reduced to a set of first-order equations:

a. Ex 1 $\frac{d^2y}{dx^2} + q(x) \frac{dy}{dx} = r(x)$

$\Rightarrow \begin{cases} \frac{dy}{dx} = z(x) \\ \frac{dz}{dx} = r(x) - q(x)z(x) \end{cases} \quad \leftarrow \text{Define new variable } z(x)$

b. The general form, then, is for a set of N ~~but~~ coupled first-order differential equations for functions

$$y_i, \quad i=1, 2, \dots, N$$

where $\frac{dy_i(x)}{dx} = F_i(x, y_1, y_2, \dots, y_N)$

and the functions F_i are known.

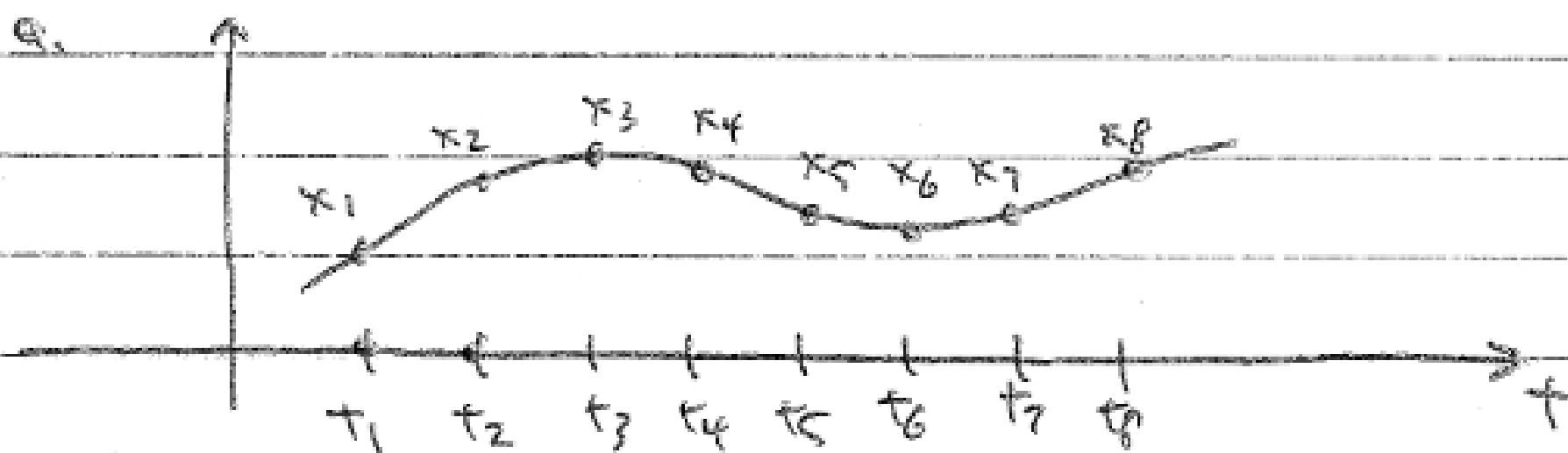
c. Boundary values on functions y_i are required to specify solution fully.

d. Plasma simulation usually requires $y_i(t=0)$ as initial conditions.

L (Continued)

B. Euler's Method:

1. To solve an equation numerically, we must represent it discretely. Take $\frac{dx}{dt} = v$



b. Take $t_j \equiv j \Delta t$ and $x(t_j) \equiv x_j$ [and $v_j \equiv v(t_j)$]

2. Discretize

$$\frac{dx}{dt} = \frac{x_{j+1} - x_j}{\Delta t} = v_j$$

a. Thus

$$x_{j+1} = x_j + \Delta t v_j$$

3. Order of Convergence

a. Take $\frac{dy}{dt} = \frac{y_{j+1} - y_j}{\Delta t} \Rightarrow y_{j+1} = y_j + \Delta t \frac{dy}{dt} \Big|_{t_j}$

b. The true solution $y(t_{j+1}) = y(t_j) + \frac{\Delta t}{1!} \frac{dy}{dt} \Big|_{t_j} + \frac{\Delta t^2}{2!} \frac{d^2y}{dt^2} \Big|_{t_j} + \frac{\Delta t^3}{3!} \frac{d^3y}{dt^3} \Big|_{t_j} + \dots$

c. Error = $y(t_{j+1}) - y_{j+1} = + \frac{\Delta t^2}{2!} \frac{d^2y}{dt^2} \Big|_{t_j} + \dots = \mathcal{O}(\Delta t^2)$

d.

DEFINE: A method is order n if error scales as $\mathcal{O}(\Delta t)^{n+1}$.

Thus, Euler's Method is first order.

C. Higher Order Methods

1. Can we conceive a higher order method for $\frac{dy}{dt}$?

I.C. (Continued)

2. Take $\frac{dy}{dt} \Big|_{t_j} = \frac{y_{j+1} - y_{j-1}}{2\Delta t}$

3. In this case, $y_{j+1} = y_{j-1} + 2\Delta t \frac{dy}{dt} \Big|_{t_j}$

a. ~~Since~~ Take $y_{j+1} = y(t_j) + \Delta t \frac{dy}{dt} \Big|_{t_j} + \frac{\Delta t^2}{2!} \frac{d^2y}{dt^2} \Big|_{t_j} + \frac{\Delta t^3}{3!} \frac{d^3y}{dt^3} \Big|_{t_j} + \dots$

b Thus $y(t_{j+1}) - y_{j+1} =$

$$y_j + \Delta t \frac{dy}{dt} \Big|_{t_j} + \frac{\Delta t^2}{2} \frac{d^2y}{dt^2} \Big|_{t_j} + \frac{\Delta t^3}{3!} \frac{d^3y}{dt^3} \Big|_{t_j} - \left[y_j - \Delta t \frac{dy}{dt} \Big|_{t_j} + \frac{\Delta t^2}{2!} \frac{d^2y}{dt^2} \Big|_{t_j} - \frac{\Delta t^3}{3!} \frac{d^3y}{dt^3} \Big|_{t_j} + \dots \right]$$

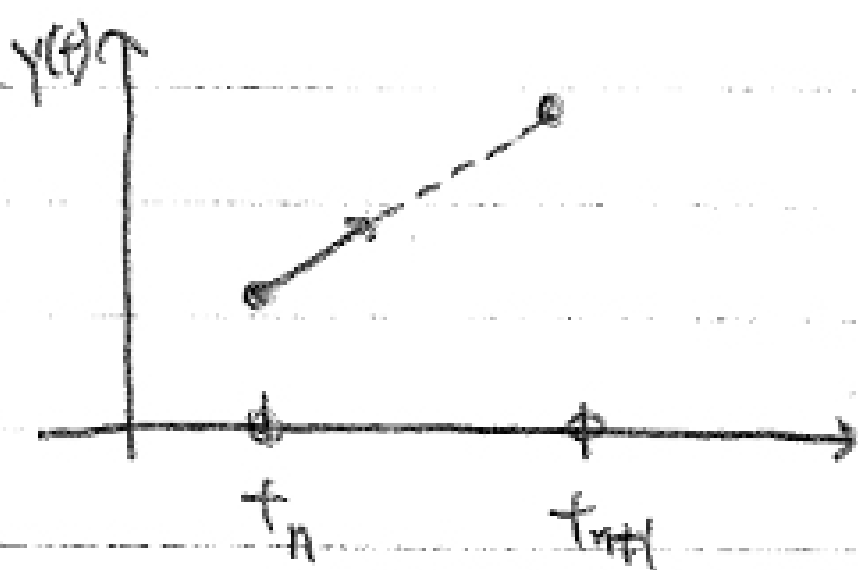
$$- 2\Delta t \frac{dy}{dt} \Big|_{t_j} = \frac{2\Delta t^3}{3!} \frac{d^3y}{dt^3} \Big|_{t_j} + \dots = O(\Delta t^3)$$

c. Thus, this centered time difference is Second order

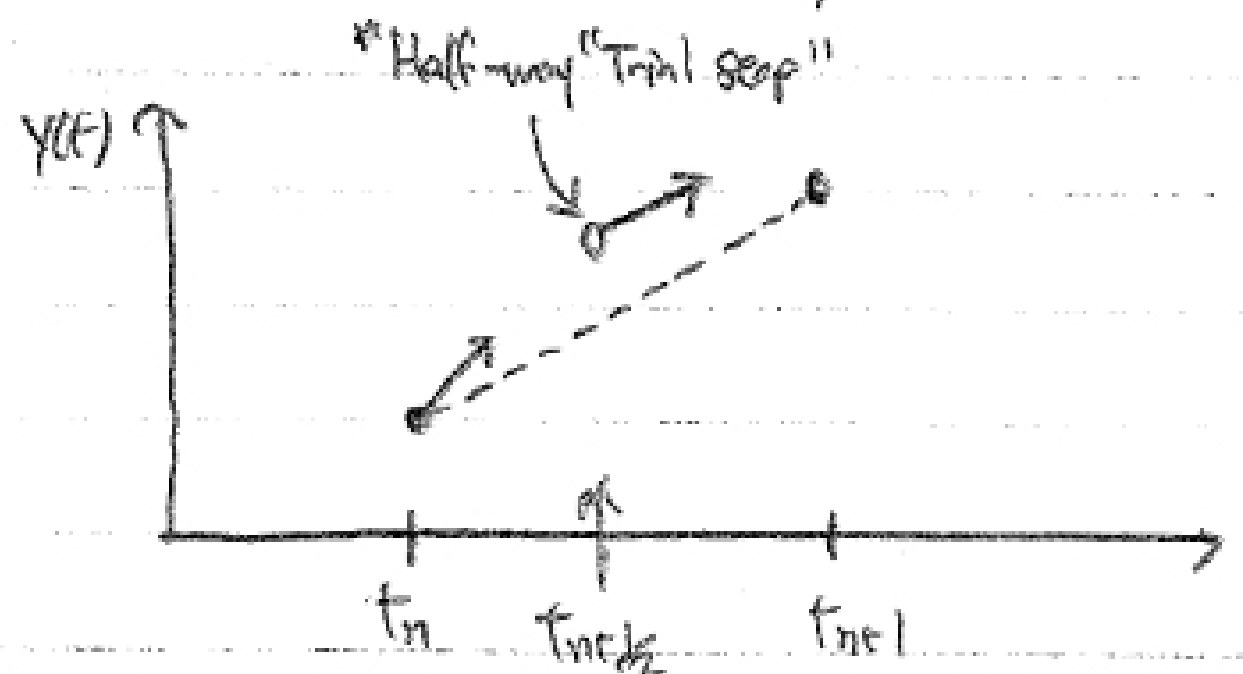
D. Runge-Kutta Methods

1. Runge-Kutta Methods are a widely used workhorse of numerical methods.

2. In this multi-step method, a "trial step" is taken to determine the derivative at the midpoint of the interval.



Euler's Method



2nd-order Runge-Kutta (or Midpoint) Method

$\frac{dy}{dt} = f \Rightarrow$ be written:

$$k_1 = \Delta t f(t_n, y_n)$$

$$k_2 = \Delta t f(t_{n+1/2}, y_n + \frac{1}{2}k_1)$$

$$y_{n+1} = y_n + k_2 + O(\Delta t^3)$$