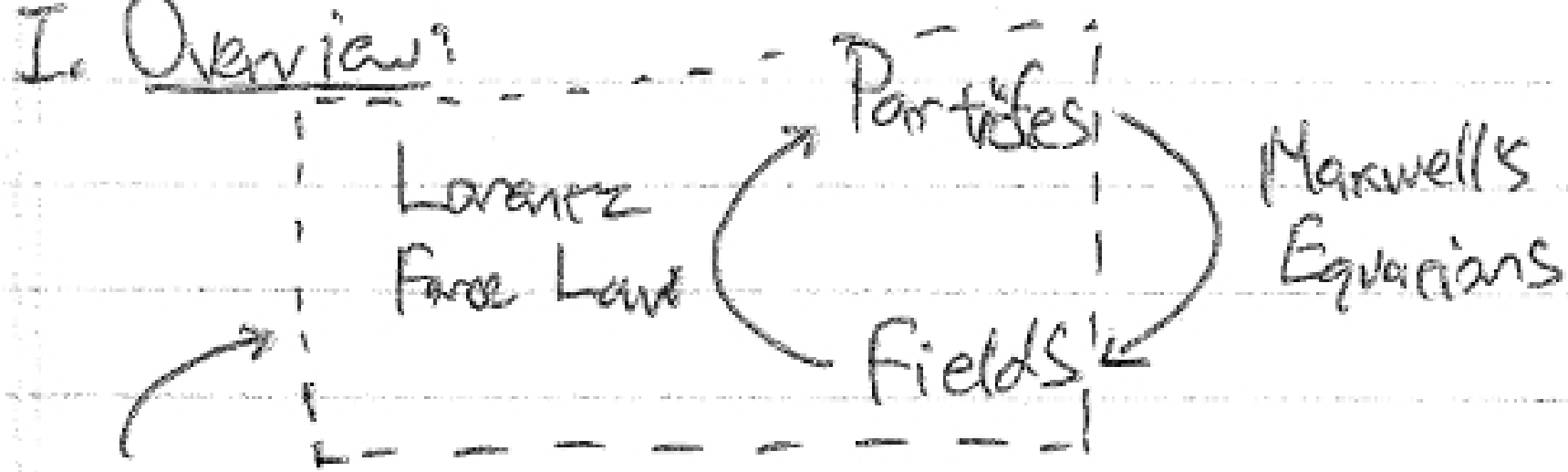


# Lecture #3: Single Particle Motion in Constant Uniform Fields Hwvcs ①

## I. Overview:



Today, and for next few weeks, we will focus on "half" of the plasma physics problem, the inconsistent picture of the effect of fields on the motion of particles, neglected the particles effect on fields.

## II. Constant, Uniform $\underline{B}$ with $\underline{E} = 0$ .

### A. 1. Nonrelativistic Limit

2. Lorentz force law  $m_s \frac{d\underline{v}_s}{dt} = q_s (\underline{E} + \underline{v}_s \times \underline{B})$

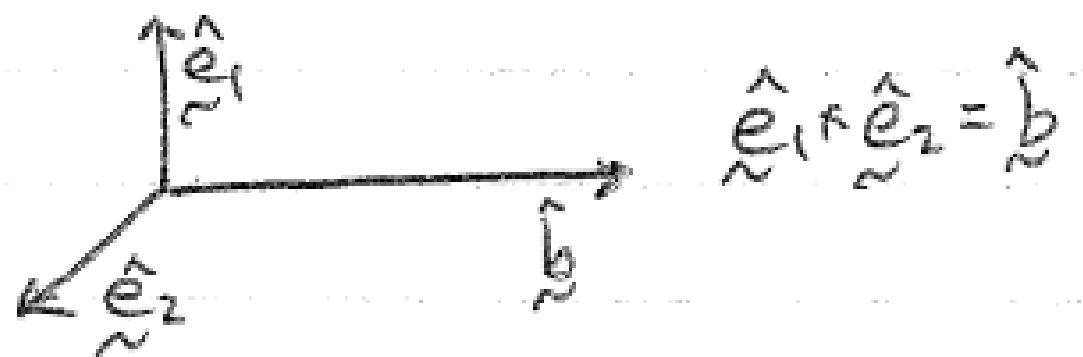
3. Drop species subscript "s" because single particle motion does not depend on other species (or any other particles),

4. Thus, for  $\underline{E} = 0$ ,  $m \frac{d\underline{v}}{dt} = q \underline{v} \times \underline{B}$

### B. Solution

1. Define  $\hat{\underline{b}} \equiv \frac{\underline{B}}{|\underline{B}|}$  and unit vectors  $\hat{\underline{e}}_1$  and  $\hat{\underline{e}}_2$  to define a

right-handed, orthonormal basis



2. We can write  $\underline{v} = v_1 \hat{\underline{e}}_1 + v_2 \hat{\underline{e}}_2 + v_{||} \hat{\underline{b}}$

and  $\underline{B} = B_0 \hat{\underline{b}}$

Thus  $\frac{d\underline{v}}{dt} = \frac{q B_0}{m} \underline{v} \times \hat{\underline{b}}$

C. NOTE:  $\underline{v} \times \hat{\underline{b}} = -v_1 \hat{\underline{e}}_2 + v_2 \hat{\underline{e}}_1$

## Lecture K3 (Continued)

Howes ③

### II. B. (Continued)

3. Take dot product with each unit vector to get component equations

$$\hat{e}_1 \cdot \frac{d\vec{v}}{dt} = \frac{qB_0}{m} \vec{v} \times \hat{b} \Rightarrow \frac{dv_1}{dt} = \frac{qB_0}{m} v_2$$

$$\hat{e}_2 \cdot \quad \quad \quad \quad \quad \quad \quad \quad \quad \Rightarrow \frac{dv_2}{dt} = -\frac{qB_0}{m} v_1$$

$$\hat{b} \cdot \quad \quad \quad \quad \quad \quad \quad \quad \quad \Rightarrow \frac{dv_{||}}{dt} = 0$$

4. Take  $\frac{d}{dt}$  of  $\hat{e}_1$  eq.  $\frac{d^2 v_1}{dt^2} = \frac{qB_0}{m} \left( \frac{dv_2}{dt} \right) = -\left( \frac{qB_0}{m} \right)^2 v_1$  substitute  $\hat{e}_2$  eq.

This is just the equation of a simple harmonic oscillator.

We define the Cyclotron Frequency  $\omega_c \equiv \frac{qB_0}{m}$

Thus,  $\frac{d^2 v_1}{dt^2} = -\omega_c^2 v_1$

5. Other properties: a) Take  $\vec{v} \cdot \frac{d\vec{v}}{dt} = 0$  with equation:

$$\vec{v} \cdot \frac{d\vec{v}}{dt} = \frac{1}{2} \frac{d v^2}{dt} = 0$$

b. So  $v^2 = \sqrt{v_1^2 + v_2^2 + v_{||}^2}$  is constant

c. But, we also know from  $\hat{b}$  equation that  $\frac{dv_{||}}{dt} = 0$ , so

$v_{\perp} \equiv \sqrt{v_1^2 + v_2^2}$  is constant.

d. Acceleration is perpendicular to  $\vec{v}$ , so Lorentz Force cannot change the energy of the particles.

6. Complex Solution: a. General Solution  $v_1 = A e^{-i\omega_c t} + B e^{i\omega_c t}$

b.  $\hat{e}_1$  component gives  $v_2 = i(-A e^{-i\omega_c t} + B e^{i\omega_c t})$

c. Take initial conditions  $v_{10} = v_{\perp}$  at time  $t=0$ .  
 $v_{20} = 0$

Lecture #3 (Continued)  
 II.B.G. (Continued)

Notes ③

d. This gives  $A=B=\frac{v_{\perp}}{2}$

e. Solution for velocity:  $v_1 = \frac{v_{\perp}}{2}(e^{i\omega t} + e^{-i\omega t}) = v_{\perp} \cos \omega t$   
 $v_2 = \frac{i v_{\perp}}{2}(e^{i\omega t} - e^{-i\omega t}) = v_{\perp} \sin \omega t$   
 $v_{\parallel} = v_{\parallel 0}$

7. Solve for Position  $x$ .

a.  $v_1 = \frac{dx_1}{dt}$  so  $x_1 = \frac{v_{\perp}}{\omega} \sin \omega t + x_{10}$

$v_2 = \frac{dx_2}{dt}$   $x_2 = \frac{v_{\perp}}{\omega} \cos \omega t + x_{20}$

$v_{\parallel} = \frac{dx_{\parallel}}{dt}$   $x_{\parallel} = v_{\parallel 0} t + x_{\parallel 0}$

b. We can define the Larmor radius (note this is not the thermal Larmor radius defined using  $v_e$ )

$$r_L \equiv \frac{v_{\perp}}{\omega} = \frac{m v_{\perp}}{q B_0}$$

c. The perpendicular components  $x_1$  &  $x_2$ , can then be written.

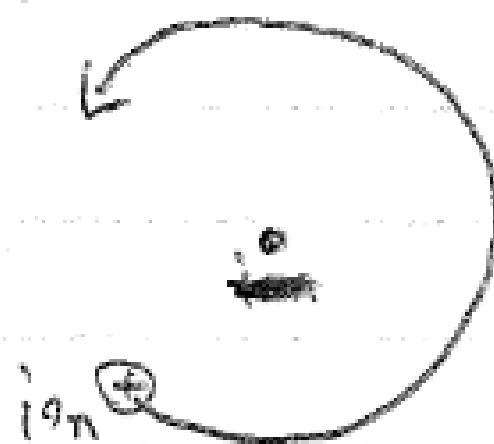
$x_1 = r_L \sin \omega t + x_{10}$

$x_2 = r_L \cos \omega t + x_{20}$

d. Larmor Motion



Since the field created by Larmor Motion opposes the mean field



⇒ DIA MAGNETIC behavior